For this preliminary experiment we have chosen to measure a 2.8% Cr-doped (V_{1-x}Cr_x)O_3 (x=0.028) single crystal. It was grown using skull melter by the Chemistry Department of Purdue University (Indiana, USA). The nominal stoichiometry was tested measuring the DC susceptibility at Laboratoire L. Neel (Grenoble). The effect of Cr-substitution at V site is to increase the Mott metal-insulator transition (PM->AFI) at T_N=180K, whereas the underdoped V sample has a reduced T_N=150K. The understanding of physical properties of these compounds at different dilution is a fundamental question addressed by our project.

V_2O_3 crystallises in the a-corundum structure (space group No. 165, R3c) and can be described by the hexagonal space group with the lattice parameters a_H=b_H=5.004 Å, c_H=13.932 Å for 2.8% Cr-doped. At T_N, the samples become antiferromagnetic with an ordered moment on the V-site of 1.2 μ_B and a propagation vector of type (1/2,1/2,0)_H. The distorted structure can be described by a monoclinic lattice (12/a) with a_m=7.277 Å, b_m=4.997 Å, c_m=5.540 Å and β=96.74° [2-5]. The cutting operations as well as the surface preparation have been performed at SPSMS/DRFMC/MDN at CEN Grenoble. The final surface etching was made at Chemical Laboratory of ESRF using a dilution of Aqua Regia. The crystal of dimension 1.0x0.5x0.05 mm³, with the planar surface orientation choose is the [110]_H, was glued with wax on a sample holder and put inside a close cycle refrigerator. The experiment was carried out tuning the incident photon energy around the K-edge of Vanadium (at E=5476 eV, λ=2.264 Å). The preliminary room temperature measurements were undertaken in order to understand the lattice structure, the quality of the surface and the absorption of the crystal. The fluorescence measured with a grazing take-off angle of 3° (Fig. 1, upper panel) shows the transitions allowed at the absorption K-edge of Vanadium. The first derivative allows us to assign the correct energies at the different states involving the dipole and quadrupole transitions. As already pointed out [6], the pre-edge 1s->3d transitions, forbidden to zeroth order, become allowed because of the presence of a small admixture of 4p-antibonding states due to the lack of inversion centre in the V-O coordinations [7]. In the AFI region the monoclinic distortion gives rise to a splitting of charge Bragg peaks in three domains, but because the thickness of the crystal used (≤50 μm), only one domain was formed. This result have permitted us to select only one magnetic domain, with a propagation vector of type (0,1,0)_m. The lower two panels of Fig. 1 shows the energy-scans taken at 100K of two magnetic reflections (221)*_m and (131)*_m, with the polarisation analysis configured (σ-π and σ-σ channels). The sharp peak at E=5465 eV corresponds to the transition associated to the 1s->3d antibonding states with t_2g(π) symmetry, in addition to a complex multiplet structure centred at E=5476 eV, observable only in the σ-π channel (central panel, absent in the lower panel), and covering the energy region where the 1s->4p transitions are allowed. The temperature dependence of the integrated intensities of the magnetic peak (221)*, at these two energies is show in the lower panel of Fig.2. Obviously, the magnetic signal disappears at T_N=180K, as determined by the DC susceptibility. Below T*_N=150K a bump appears which is more intense in the quadrupole transition.
associated to the 3d-t_{2g}(\pi) states. At this temperature the lattice constant, determined by the longitudinal scans around the (231)^* charge peak (central panel of Fig 2), show an anomalous behaviour in a continuous manner, in contrast with what is observed in pure V_2O_3. Note that 150K corresponds to the metal-insulator magnetic transition in the pure compound. The last part of the experiment was dedicated to the search of the peaks of the orbital ordering. In fact, as suggested by many authors [8-9] and following the calculations and suggestions of the theoretical group at ESRF [10], the orbital order in V_2O_3 is expected to break the translation symmetries giving rise to a propagation vector located at different positions in reciprocal lattice with respect to the antiferromagnetic one. The two different orbital structures proposed correspond to propagation vectors ((1/3,1,-1/3)^*, and (2/3,1,1/3)^*, these reciprocal vectors are expressed in our monoclinic domain, and are equivalent to the hexagonal reflections ((1,-1/2,1)^*H and (2/3,1,1/3)^*H, respectively). Despite the lack of orbital order peaks at the positions proposed, these results are important and motivate the continuation of this project. The main question arising concerns the possibility to have a different behaviour in the 2.8% Cr-doped V_2O_3 single crystal with respect to the pure compounds. In fact, from our results it is tempting to ascribe the T^* =150K to the orbital ordering temperature. In this case leading to a two different and separate order parameters where the magnetism drive the orbital order. This fact can explain why the Cr-doped samples have a modified structural transition with respect to the pure. This assumption is confirmed by inelastic neutron scattering [11] where just below the TN a big increase of inelastic signal is present in samples with different dilution. The bump exist also in the recent magnetoresistance measurements and it is suppressed by the pressure [12].

References

![Fig. 1](image1)  
**Fig. 1** Energy scans of Cr-doped sample. The panel 1-a shows the K-photon absorption spectrum (filled points) taken at room temperature. The capital letters mark the different transitions determined from the maxima of the absorption first derivative (wide points): the peak B corresponds to the transitions to 3d anti-bonding states with symmetries t_{2g}(\pi) [see Ref. 6]. The two panels b-c show the E-scans of antiferromagnetic Bragg peaks [221]_m and [131]_m taken at T=100K in \sigma-\sigma (wide points) and \sigma-\pi (black points) polarisation analysis.

![Fig. 2](image2)  
**Fig. 2**  
a) DC susceptibility showing the Néel magnetic transition T_N =180K (doped sample). The filled (wide) point is the heating (cooling) measurements.  
b) Temperature evolution of the centre position of charge peak [231]_m taken by a scan along the longitudinal direction \langle H, H, 0 \rangle_{\text{mon}}. T^* =150K is the lattice distortion temperature from rhombohedral to monoclinic structure  
c) Temperature dependence of the integrated intensities of the antiferromagnetic Bragg peak [221]_m at the dipole and at the quadrupole resonant energies of the Vanadium K-edge. The bump appearing at T^* is enhanced at the quadrupole energy.