



	Experiment title: Two-dimensional x-ray waveguide structures	Experiment number: Mi439
Beamline: ID 13	Date of experiment: from: 22.11.2000 to: 29.11.2000	Date of report: 06.03.2001
Shifts: 21	Local contact(s): C. Riekkel, M. Burghammer	<i>Received at ESRF:</i>

Names and affiliations of applicants (* indicates experimentalists):

Dr. Tim Salditt* – Universitaet des Saarlands, Experimentalphysik, Im Stadtwald 38, Postfach 15 11 50, 66041 Saarbruecken, Germany, phone: +49 681 302 2216, email: salditt@mx.uni-saarland.de

Franz Pfeiffer* – Sektion Physik der Universität München, Geschwister-Scholl-Platz 1, D-80539 München, Germany, phone: +49 681 302 3032, email: franz.pfeiffer@physik.uni-muenchen.de

Dr. Christian David* – Laboratory for Micro- and Nanotechnology, Paul Scherrer Institut, CH-5232 Villigen-PSI, phone: +41 56 310 3753, email: christian.david@psi.ch

Report:

Planar x-ray waveguide structures as resonant beam coupling devices have been demonstrated as new tools to produce coherent and divergent x-ray beams with cross-sections in the sub-micrometer range in one dimension. However, for many nano-microscopy and nano-spectroscopy applications, a two-dimensionally confined *point* beam instead of one-dimensional *line* beam is needed.

In this experiment we have demonstrated for the first time that a two-dimensional beam compression using the resonant beam coupling principle can be achieved. Applying state of the art e-beam lithography techniques, we have fabricated laterally and vertically patterned nanostructures, where the beam is channeled in two-dimensions. Figure 1 shows a schematic of such a 'x-ray fiber' consisting of an approximately $500\text{\AA} \times 2000\text{\AA}$, spin-coated, rectangular patterned PMMA core coated with evaporated Cr metal on silicon substrate.

Just like in the optical analogue, the resonantly enhanced electromagnetic field distribution of a x-ray waveguide can be described by a set of discrete and precisely defined modes (at the respective angles α_i and ϕ_i), which are determined by the geometry of the structure and the choice of materials. Accordingly, the guided modes are excited by shining a parallel beam onto the waveguide at grazing incidence at a set of angles $\alpha_{i,n}$ and $\phi_{i,m}$. A coherent beam

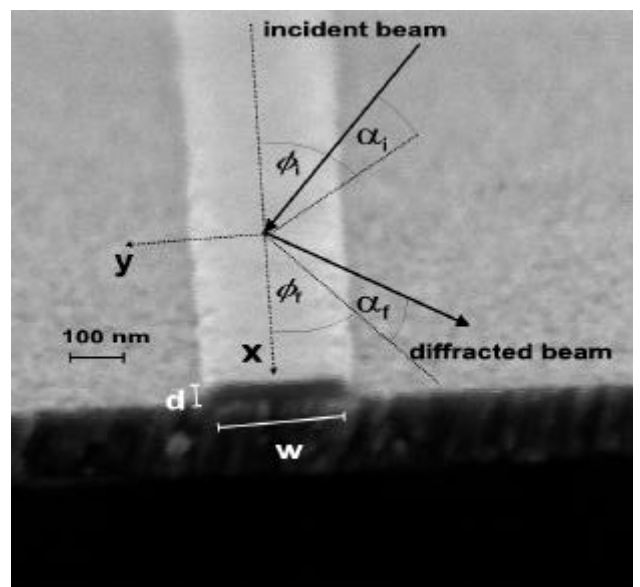


Fig. 1: Scanning electron micrograph of a two-dimensional waveguide consisting of an e-beam patterned PMMA-core coated with Cr.

exists the structure at the side with a cross section corresponding to the thickness of the guiding layer and a divergence in the range of some mrad.

To experimentally demonstrate the waveguide effect due to resonant beam coupling, we have measured the reflectivity of the nanostructure both as a function of α_i ($\phi_i = \text{const.}$, $\alpha_f = 2\alpha_i$) and ϕ_i ($\alpha_i = \text{const.}$, $\phi_f = 2\phi_i$). The waveguide effects readily manifest themselves in the reflectivity curve as cusps in the plateau of otherwise

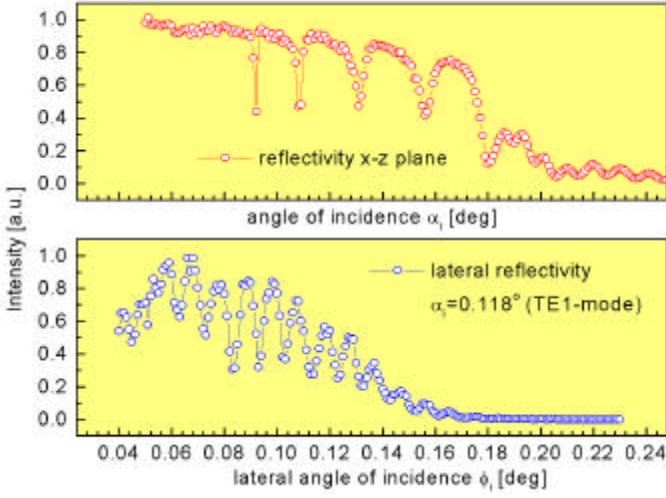


Fig. 2: Reflectivity of the waveguide nanostructure both as a function of α_i ($\mathbf{f}_i = \text{const.}$, $\alpha_f = 2\alpha_i$, top) and ϕ_i ($\alpha_i = \text{const.}$, $\mathbf{f}_f = 2\mathbf{f}_i$, bottom).

total external reflection. Importantly, the evidence for the resonances can be clearly seen in both the α_i / α_f [Fig. 2 top] and ϕ_i / ϕ_f – geometry [Fig. 2 bottom, respectively]. In a simplistic argument valid for infinite samples and beams, photons get trapped under the resonance conditions in the guiding layer propagating over an active coupling length (in our case $\sim 500\mu\text{m} - 1000\mu\text{m}$) parallel to the tunnel, and are therefore more likely to get absorbed. The enhanced absorption loss manifests itself as a pronounced decrease in the reflectivity. At the same time, the ‘point source’ has been characterized by translating a knife edge in a distance of $\sim 1.8\text{mm}$ across the divergent beam exiting the end of the ‘fiber’ for the $\alpha_{i,0} / \phi_{i,0}$ – resonance [Fig. 3]. The resulting FWHM – beam crosssections of $4.1\mu\text{m} / 8.4\mu\text{m}$ in z-/y-direction correspond well with the expected value of $3.6\mu\text{m}$ in that distance due to the $\sim 0.12^\circ$ divergence of the exiting mode. The broadening (particularly in y-direction) is most likely caused by the diffuse background of the primary beam and diffraction at the imperfect knife-edge.

To gain an more detailed understanding of the properties of the resonance-excitation of modes in such patterned structure, their farfield intensity-distributions, the achieved waveguide gain and the absolute performance of the device a series of detector mappings as a function α_i and ϕ_i were recorded with a high resolution CCD camera. The corresponding data evaluation is still in progress, Fig. 4 already gives a first idea of the complexity of the diffraction pattern and thus of the potential of twodimensional x-ray waveguides. More detailed information, particularly concerning the topics discussed above (farfield-intensity-distribution, performance, gain) as well as input for improved structures will be published within the next few weeks.

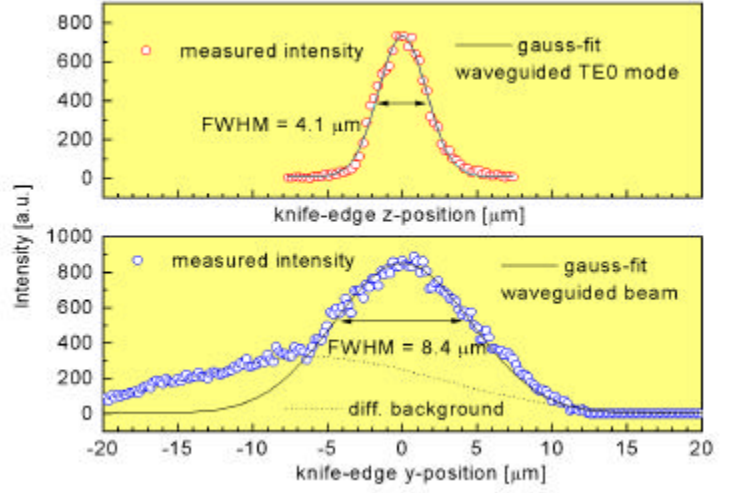


Fig. 4: Beam-crosssections in a distance of $\sim 1.8\text{mm}$ across the divergent beam exiting the end of the ‘fiber’ for the $\alpha_{i,0} / \phi_{i,0}$ – resonance

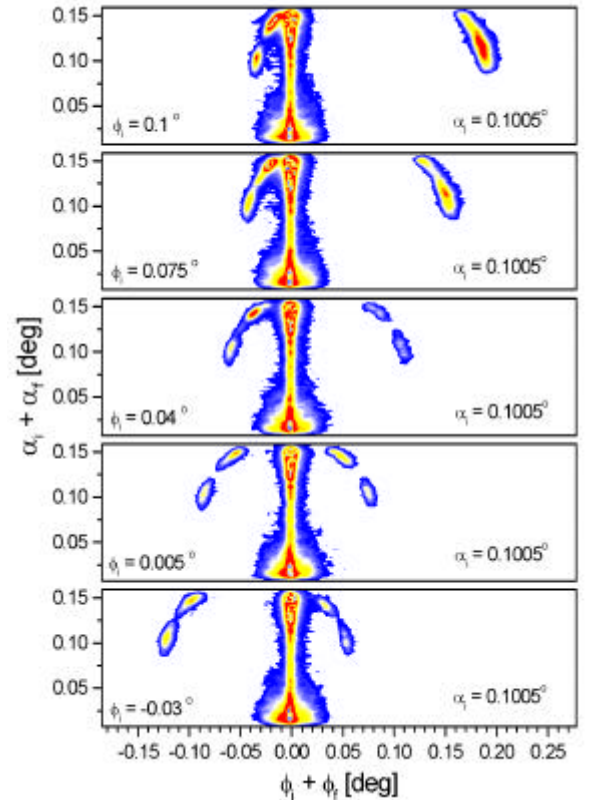


Fig. 4: Contour-plots of a series of detector mappings collected for various α_i / ϕ_i values.