

**Experiment title:**

Diffractometric determination of the interface stress in ferro- and paramagnetic nanocrystalline Gd

**Experiment number:**

SI-788

**Beamline:**

ID31

**Date of experiment:**

from: 01 July 2002, 7:00 to: 02 July 2002, 7:00

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3

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**Report:**

The goal of Experiment Si-788 was to obtain a value for the area-averaged interface (i.e. grain-boundary) stress  $\langle f \rangle_A$  in nanocrystalline Gd from a measurement of the volume-averaged (and, therefore, residual-strain-free) lattice parameters. A previous, indirect measurement of  $\langle f \rangle_A$  suggested an anomalously large value for the interface stress, offering an explanation for the measured dependence of the Curie transition temperature  $T_C$  of Gd on the average crystallite size [1]. Because  $\langle f \rangle_A$  induces a change  $\Delta v$  in the volume of the unit cell  $v_0$  according to

$$\frac{\Delta v}{v_0} = -\frac{4}{3} \frac{\langle f \rangle_A}{B^*} \left( \frac{1}{L_{\text{area}}} \right), \quad (1)$$

where  $B^*$  denotes the (polycrystalline) bulk modulus and  $L_{\text{area}}$  the area-weighted average crystallite size, it is possible to extract the average interface stress from a careful measurement of lattice parameters as a function of  $L_{\text{area}}$ . Since such measurements can be performed both below and above the Curie temperature, we hoped not only to obtain a direct measure of the true value of  $\langle f \rangle_A$  in our samples, but also to determine whether ferromagnetic ordering (and the resulting magnetostriction) exerts any influence on the interface stress.

Samples of nanocrystalline Gd were prepared by inert-gas condensation/compaction in the form of disks ( $\phi$  8 mm) having a thickness of 0.3–0.5 mm. The disks were cut into narrow strips and sealed in glass capillaries ( $\phi$  0.7 mm), which were then mounted in the high-resolution powder diffractometer of beamline ID31. A streaming-LN<sub>2</sub> cryostat was employed to perform *in situ* measurements below room temperature, whereas elevated temperatures were reached with the help of a hot-air blower mounted just below the rotating capillary. Wide-angle diffraction scans were recorded from  $-4^\circ$  to  $16^\circ$  ( $2\theta$ ) with x-ray photons of  $\lambda = 0.24900 \text{ \AA}$  ( $E = 49.79 \text{ keV}$ ) while moving the nine-channel detector at  $1^\circ/\text{min}$ . The instrumental broadening at this wavelength was determined by fitting the profiles of Bragg peaks produced by a coarse-grained standard sample (Na<sub>2</sub>Ca<sub>3</sub>Al<sub>2</sub>F<sub>14</sub>) provided by the beamline.

Seven samples with average crystallite sizes ranging from  $L_{\text{area}} = 4 \text{ nm}$  to  $65 \text{ nm}$  were investigated at 100 K, room temperature ( $\sim 295 \text{ K}$ ) and  $70^\circ\text{C}$  (343 K). Since the Curie temperature of coarse-grained Gd is 292 K and the  $T_C$  of Gd is known to decrease with  $L_{\text{area}}$ , reaching 266 K at 5 nm [1], we expect the samples to be ferromagnetic at 100 K and paramagnetic at the two higher temperatures.

According to Eq. (1), we can determine the area-averaged interface stress  $\langle f \rangle_A$  from measurements of the relative volume change of the unit cell,  $\Delta v/v_0$ , as a function of the area-weighted average crystallite size  $L_{\text{area}}$ . The latter can be extracted from an analysis of Bragg-peak profiles [2], while the volume change can be obtained at a given temperature by comparing the values of the lattice parameters  $a$  and  $c$  of the hexagonal unit cell to the literature values for single-crystalline Gd ( $a = 363.6$  pm,  $c = 578.26$  pm at room temperature [3]); values for  $v_0$  at 100 K and 343 K were estimated using thermal-expansion data for polycrystalline Gd [4]. We determined both the lattice parameters and the crystallite size by performing a full-pattern profile fit of each diffraction scan using the commercially available program FORMFIT [5]. This program performs a Rietveld-like refinement of (i) the parameters of the pseudo-Voigt function that is used to approximate the profiles of individual Bragg peaks and (ii) the lattice parameters, which fix the centers of the individual peak fits. The peak-profile function is convoluted with the instrumental broadening, with the microstrain and with an assumed lognormal crystallite-size distribution, thus permitting simultaneous refinement of the lattice parameters  $a$  and  $c$ , the volume-weighted crystallite size  $L_{\text{vol}}$ , the standard deviation  $\sigma$  of the lognormal distribution and the microstrain  $e$ . From the values of  $L_{\text{vol}}$  and  $\sigma$  we calculated the area-weighted crystallite size  $L_{\text{area}}$  according to  $L_{\text{area}} = [(8/9) \exp(-\ln^2 \sigma)]L_{\text{vol}}$ , an expression that is valid for a lognormal distribution of spherical-shaped crystallites [2].

The experimental data plotted in Fig. 1 show no sign of a dependence of  $\Delta v/v_0$  on the magnetic ordering in the samples, as the data points measured at all three temperatures appear to lie along a common curve. The fact that this curve is exponential rather than linear, as expected for a constant value of  $\langle f \rangle_A$ , indicates either that the value of the interface stress depends on the average crystallite size, or that the samples differ from each other not only in the value of  $L_{\text{area}}$ . Given the bulk modulus  $B^*$  of polycrystalline Gd (37.9 GPa), we estimate an upper bound of  $\langle f \rangle_A \approx 2$  N/m from the slope of the fit curve at  $L_{\text{area}} = 100$  nm, whereas the slope at a crystallite size of 10 nm yields  $\langle f \rangle_A \approx 0.2$  N/m. Even the larger of these two values for the average interface stress is well below the level of 5.5 N/m needed to explain the crystallite-size dependence of  $T_C$  that is observed in nanocrystalline Gd [1]. Thus, we conclude that the Curie-temperature depression measured in nanocrystalline Gd cannot result entirely from the volume contraction induced by the interface stress of the grain boundaries. Likewise, it is probable that observations of a drastic dependence of  $T_C$  on the thickness of Gd thin films [6] cannot be explained entirely by the *surface* stress, as the latter is generally a factor of only 2 or 3 larger than  $\langle f \rangle_A$ .

[1] D. Michels, C. E. Krill III and R. Birringer, *J. Mag. Mag. Mater.* **250** (2002) 203.

[2] C. E. Krill and R. Birringer, *Phil. Mag. A* **77** (1998) 621.

[3] F. H. Spedding, A. H. Daane and K. W. Herrmann, *Acta Crystallogr.* **9** (1956) 559.

[4] *Thermal Expansion: Metallic Elements and Alloys*, edited by Y. S. Touloukian, R. K. Kirby, R. E. Taylor and P. D. Desai (IFI/Plenum, New York, 1975), p. 107.

[5] [www.anadat.com](http://www.anadat.com)

[6] M. Farle, K. Baberschke and U. Stetter, *Phys. Rev. B* **47** (1993) 11571.

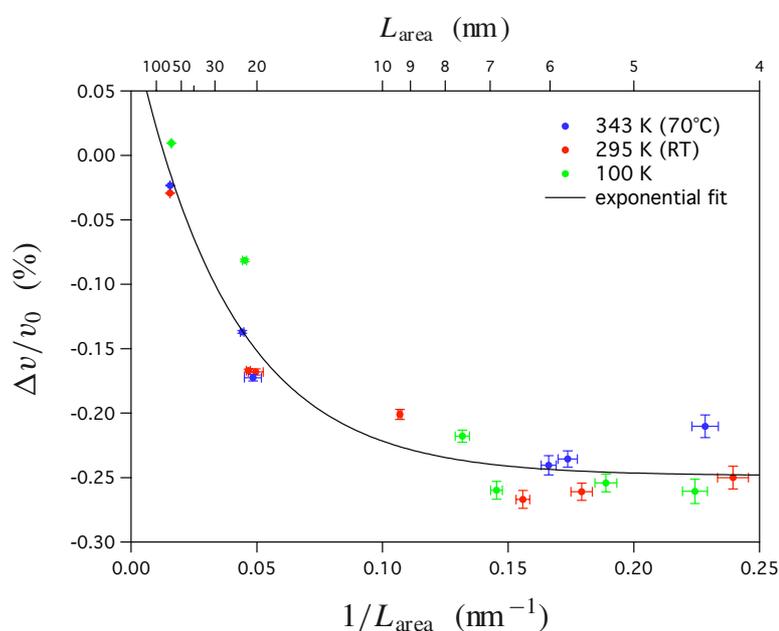


Fig. 1: Relative volume change  $\Delta v/v_0$  plotted as a function of the area-weighted average crystallite size  $L_{\text{area}}$  in nanocrystalline Gd, measured at 100 K, room temperature and 70°C.

The solid curve is a global fit of an exponential function to the data points. From the slope of the fit at  $L_{\text{area}} = 100$  nm and Eq. (1), we estimate an apparent interface stress of 2 N/m, dropping to  $\sim 0.2$  N/m at crystallite sizes below 10 nm.