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- fill in a separate form for each project or series of measurements.
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- if your work is published or is in press, you may prefer to paste in the abstract, and add full reference details. If the abstract is in a language other than English, please include an English translation.



	Experiment title: Monoclinic distortions and orbital-ordering in the Ca-doped Mott insulator $\text{Pr}_{1-x}\text{Ca}_x\text{VO}_3$	Experiment number: HS-4297
Beamline:	Date of experiment: from: 05/05/2011 to: 09/05/2011	Date of report: 27/05/2011
Shifts:	Local contact(s): P. M. Abdala	<i>Received at ESRF:</i>
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Report:

Ternary transition-metal oxides with perovskite structure have attracted considerable attention due to their unusual phase behavior and physical properties, which are consequences of an intricate interplay between charge, orbital, and spin ordering. Orthovanadates of chemical composition $R\text{VO}_3$, with $R = \text{Y}$ or a trivalent rare-earth metal atom, are particularly interesting examples, as they undergo multiple orbital and magnetic ordering transitions as a function of temperature [1-6]. At room temperature these materials exhibit the orthorhombic GdFeO_3 -type structure with space group $Pbnm$. With decreasing temperature, they undergo a structural phase transition (T_S), ascribed to orbital ordering, into a phase described by the monoclinic space group $P2_1/b$, and a Néel transition (T_N) associated with the formation of a C -type magnetic structure of the spin-1 moments of V^{3+} . For PrVO_3 these transitions set in at $T_S = 171$ K and $T_N = 140$ K, respectively. Recently the electronic phase diagram has been investigated for the Ca-doped vanadates $\text{Pr}_{1-x}\text{Ca}_x\text{VO}_3$ [4]. The increase of the doping level $x(\text{Ca})$ causes a melting of the orbital and spin order followed by a metal-insulator transition at $x(\text{Ca}) = 0.23$.

In order to investigate in detail the orthorhombic and monoclinic structure of Ca-doped vanadates $\text{Pr}_{1-x}\text{Ca}_x\text{VO}_3$ we performed high-resolution powder diffraction studies at the ESRF in Grenoble. For the experiment we took single crystal samples, and ground fragments to fine powders, which were mounted in a glass capillary of diameter 0.5 mm. Complete powder patterns on the BM1B (Swiss-Norwegian) beam line with a 2θ -range between 1 and 55° were collected using the wavelength $\lambda = 0.50214(4)$ Å. Rietveld refinements of the powder diffraction data were carried out with the program *FullProf* using the atomic scattering factors

provided by the program [7]. Refinements of the diffraction patterns PrVO_3 , $\text{Pr}_{0.90}\text{Ca}_{0.10}\text{VO}_3$ and $\text{Pr}_{0.80}\text{Ca}_{0.20}\text{VO}_3$ collected at 15 K were performed in the monoclinic space group $P2_1/b$. For comparison we collected a data set of the orthorhombic phase (space group $Pbnm$) of $\text{Pr}_{0.90}\text{Ca}_{0.10}\text{VO}_3$ at 290 K. Here the V- and O2-atoms are located in the Wyckoff positions $4b(\frac{1}{2},0,0)$ and $8d(x,y,z)$, respectively, while the O1- and Pr(Ca)-atoms are located at the position $4c(x,y,\frac{1}{4})$. In the monoclinic phase the vanadium atoms are located in the positions $2c(\frac{1}{2},0,0)$ and $2d(\frac{1}{2},0,\frac{1}{2})$, while the Pr(Ca)- and the three different O-atoms (O1, O2 and O3) are located in the position $4e(x,y,z)$. During the refinements a total of 12 and 7 positional parameters were allowed to vary in the space groups $P2_1/b$ and $Pbnm$, respectively. The crystal structure refinements of the three compounds PrVO_3 , $\text{Pr}_{0.90}\text{Ca}_{0.10}\text{VO}_3$ and $\text{Pr}_{0.80}\text{Ca}_{0.20}\text{VO}_3$ resulted in very satisfactory R_F -values between 0.015 and 0.023. The refined positional parameters finally allowed us to determine precisely the interatomic distances. The bond distances between the V- and O-atoms in the VO_6 -octahedra as well as the lattice parameters are presented in Fig. 1.

The onset of orbital ordering appears at the structural phase transition temperature T_S , while the C-type spin ordering of the V-sublattice sets in at the lower lying Néel temperature T_N . The increase of the doping level $x(\text{Ca})$ causes a decrease of the transition temperatures T_S and T_N as well as the lattice parameters a and b , while the c -axis remains almost unchanged (Fig. 1). In the range $0 \leq x(\text{Ca}) \leq 0.20$ the cell volume shows a relatively strong decrease of 1.93%. The change of V-O-bond distances in the two different VO_6 -octahedra (of V1 and V2) is shown in the lower part of the diagram of Fig. 1. Up to a doping level of $x(\text{Ca}) = 0.20$ the V1-O22- and V2-O32-bonds are strongly elongated indicating the presence of the *Jahn-Teller* effect (Fig. 1). Interestingly the *Jahn-Teller* effect is still established at the doping level $x(\text{Ca}) = 0.20$. This is close to the critical level $x(\text{Ca}) = 0.23$, where the metal-insulator transition was found [4]. The change of the apical bond lengths $d(\text{V1-O1})$ and $d(\text{V2-O1})$ is quite unusual. In PrVO_3 these bond lengths [$d(\text{V1-O1}) = 1.981(8) \text{ \AA}$, $d(\text{V2-O1}) = 1.985(8) \text{ \AA}$] are very similar. But with increasing $x(\text{Ca})$ these bond lengths differ significantly. For $\text{Pr}_{0.80}\text{Ca}_{0.20}\text{VO}_3$ finally the values $d(\text{V1-O1}) = 2.008(11) \text{ \AA}$ and $d(\text{V2-O1}) = 1.937(11) \text{ \AA}$ have been obtained.

The Ca^{2+} and Pr^{3+} ions are surrounded by 8 O^{2-} -ions in a distorted cubic coordination. All the 8 bond distances in PrVO_3 show different values between 2.367(3) and 2.713(6) \AA . A very similar range between 2.363(3) and 2.708(8) \AA could be observed for the Ca-doped vanadate $\text{Pr}_{0.80}\text{Ca}_{0.20}\text{VO}_3$. For both compounds the averaged d -values $d_{\text{av}} = 2.534(5) \text{ \AA}$ (PrVO_3) and $d_{\text{av}} = 2.535(7) \text{ \AA}$ ($\text{Pr}_{0.80}\text{Ca}_{0.20}\text{VO}_3$) are practically the same. This can be ascribed to the fact that the ionic radii of both ions Ca^{2+} ($r = 1.12 \text{ \AA}$) and Pr^{3+} ($r = 1.126 \text{ \AA}$) are very similar [8]. On the other hand an increase of the Ca-doping level [$0 \leq x(\text{Ca}) \leq 1$] changes the charge of the vanadium ions from 3+ to 4+ assuming a pure ionic system. Accordingly one expects for $\text{Pr}_{0.90}\text{Ca}_{0.10}\text{VO}_3$ and $\text{Pr}_{0.80}\text{Ca}_{0.20}\text{VO}_3$ a mixed valency with charges of +2.9 and +2.8, respectively. Due to the different ionic radii $r(\text{V}^{3+}) = 0.640 \text{ \AA}$ and $r(\text{V}^{4+}) = 0.58 \text{ \AA}$ (in an octahedral coordination) [8] the averaged d -value (from the three different bond distances in the VO_6 -octahedra) decreases with increasing Ca-doping level. In fact, we observed a slight decrease of the averaged d -value: $d_{\text{av}} = 2.005(7) \text{ \AA}$ (PrVO_3) $>$ $d_{\text{av}} = 1.994(8) \text{ \AA}$ ($\text{Pr}_{0.90}\text{Ca}_{0.10}\text{VO}_3$) $>$ $d_{\text{av}} = 1.985(9) \text{ \AA}$ ($\text{Pr}_{0.80}\text{Ca}_{0.20}\text{VO}_3$). At 290 K we obtained for $\text{Pr}_{0.90}\text{Ca}_{0.10}\text{VO}_3$ the value $d_{\text{av}} = 1.9935(14)$. Despite to the fact that the *Jahn-Teller* effect is only active in the monoclinic phase the d_{av} -values are almost unchanged.

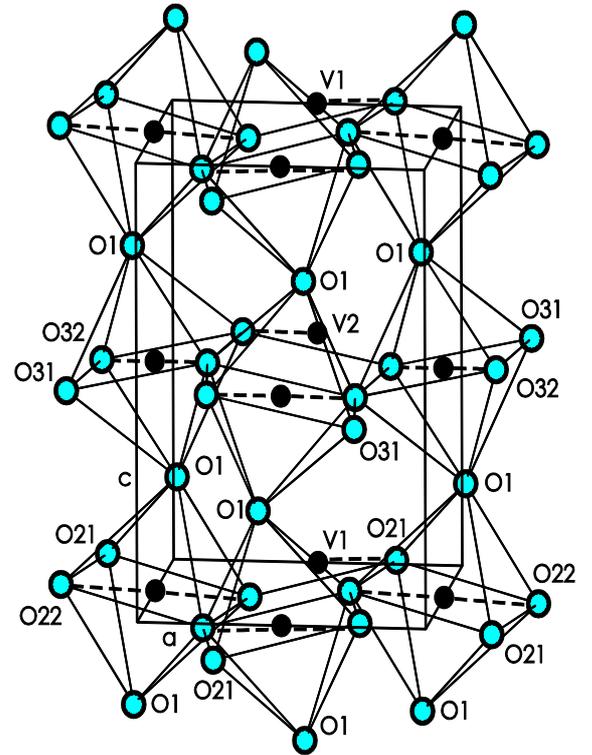
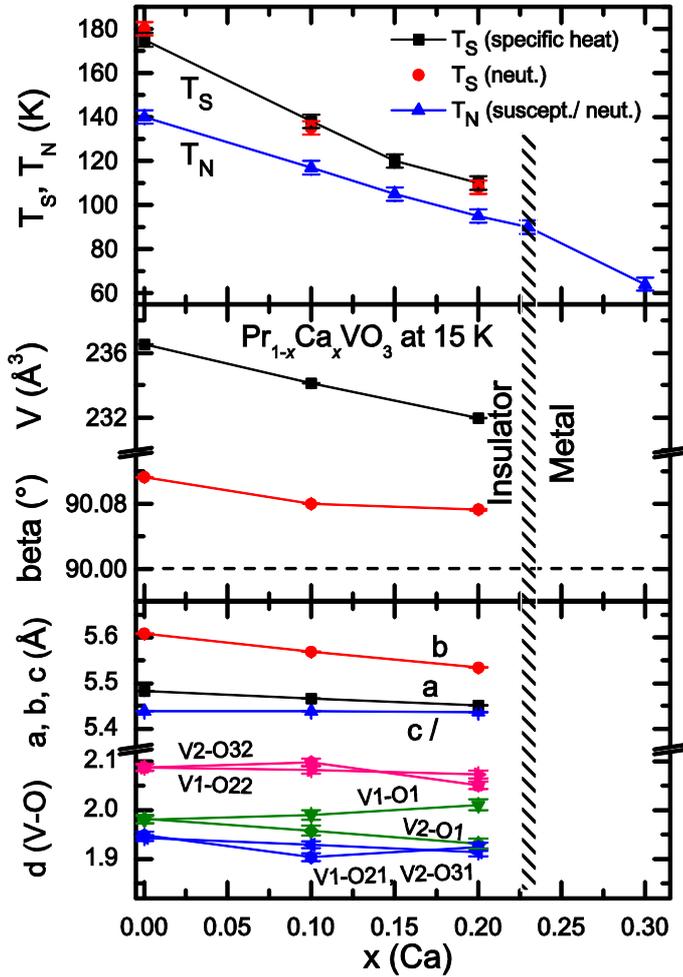


Fig. 1. Change of structural and magnetic properties in the system $\text{Pr}_{1-x}\text{Ca}_x\text{VO}_3$ as a function of the doping level $x(\text{Ca})$. In the insulating state the V1-O22- and V2-O32-bonds are strongly elongated due to the presence of the *Jahn-Teller* effect. The network of distorted corner-shared VO_6 -octahedra in the monoclinic phase of $\text{Pr}_{1-x}\text{Ca}_x\text{VO}_3$ is also shown. The dashed lines represent the elongated V-O-bonds due to the *Jahn-Teller* effect.

References

1. G. R. Blake, T. T. M. Palstra, Y. Ren, A. A. Nugroho, A. A. Menovsky, Phys. Rev. Lett. **87**, 245501 (2001); Phys. Rev. B **65** (2002) 174112.
2. S. Miyasaka, Y. Okimoto, M. Iwama, and Y. Tokura, Phys. Rev. B **68** (2003) 100406(R).
3. C. Ulrich, G. Khaliullin, J. Sirker, M. Reehuis, M. Ohl, S. Miyasaka, Y. Tokura, B. Keimer, Phys. Rev. Lett. **91** (2003) 257202.
4. J. Fujijoka, S. Miyasaka, Y. Tokura, Phys. Rev. B **72** (2005) 024460.
5. M. Reehuis, C. Ulrich, P. Pattison, B. Ouladdiaf, M. C. Rheinstädter, M. Ohl, L. P. Regnault, S. Miyasaka, Y. Tokura, B. Keimer, Phys. Rev. B **73** (2006) 094440.
6. M. Reehuis, C. Ulrich, P. Pattison, S. Miyasaka, Y. Tokura, B. Keimer, Phys. Rev. B **73** (2006) 094440.
7. J. Rodriguez-Carvajal, Physica B **192**, 55 (1993).
8. R. D. Shannon, Acta Cryst. A **32** (1976) 751.