ESRF	Experiment title: XMCD study of the behavior of Co orbital magnetization of the La <sub>1-x</sub> Sr <sub>x</sub> Co <sub>1-y</sub> Nb <sub>y</sub> O <sub>3</sub> single crystals in the range of spin-state and ferromagnetic phase transitions	Experiment number: HE-3677
Beamline: ID12	Date of experiment:   from: 31.08.2011   to: 06.09.2011	<b>Date of report</b> : 01.09.2013
Shifts: 18	Local contact(s): Dr. Andrei ROGALEV	Received at ESRF:
Dr. Vadim SIKOLENKO - Karlsruhe Institute of Technology, Germany		
Dr. Alex IGNATOV*- Materials Science & Engineering, Rutgers, NJ, USA		
Prof. Igor TROYANCHUK*- Scientific-Practical Materials Research Centre of NAS, Belarus		
Vadim EFIMOV*- Joint Institute for Nuclear Research, Dubna, Moscow region, Russia		

A perovskite-like lanthanum cobalt oxide LaCoO<sub>3</sub> belongs to a family of perovskites that attracts a lot of interest because of their rich variety of physical properties. In the ground state the Co<sup>3+</sup> ions are found to be in the low-spin state (LS;  $t_{2g}^6 e_g^0$ , S = 0). The transition ~120 K was ascribed to the thermal activation of the intermediate-spin state (IS;  $t_{2g}^5 e_g^1 e_g^1$ , S = 1) or high-spin state (HS;  $t_{2g}^4 e_g^2$ , S = 2) of Co<sup>3+</sup> [1-9].

Soft X-ray absorption spectroscopy (XAS) and X-ray magnetic circular dichroism (XMCD) at the Co  $L_{2,3}$ edges well as theoretical as calculations for a CoO<sub>6</sub> clusters [2] shown that the have spin-state transition in LaCoO<sub>3</sub> can be well described by a LS ground state and a triply degenerate HS excited state. The authors [2] argue that large orbital momentum revealed by the Co  $L_{2,3}$ edges XMCD measurements invalidates existing LS-IS scenarios [1].



It is worth noticing that recent Co  $L_{2,3}$  XMCD data by Burnus *et. al.* [6] showed that Co<sup>2+</sup> ions in

**Fig. 1:** Co *K*-edge XANES and temperatue-dependent XMCD spectra under external magnetic field of 17 Tesla for a  $LaCoO_3$  single crvstal parallel to the *c*-axis.

LaMn<sub>0.5</sub>Co<sub>0.5</sub>O<sub>3</sub> have a large orbital moment,  $L_z/S_z \sim 0.47$ . The  $L_z/S_z$  is surpisingly similar to one reported by Haverkort *et. al.* [2] in LaCoO<sub>3</sub> single crystal. Clearly, one need to be concirned about strain field and vacancies on the LaCoO<sub>3</sub> surface [7-9]. The present study was undertaken to provide a essentially buld probe of Co spins states, by making use of hard x-ray XMCD at the Co *K*-edge.

Figure 1 shows the temperature dependence of the normalized XANES and XMCD spectra for a LaCoO<sub>3</sub> single crystal recorded at Co *K*-edge in total fluorescence yield with electric field vector of X-ray perpendicular to the magnetic *c*-axis. A quite clear dichroic signal of the order of 0.02 % at T = 3 K with respect to the edge jump is visible at preedge feature of Co.

Co 3*d* and 4*p* orbital magnetic moments obtained from integrated XMCD singal shown in Fig. 1 following the prescription in Ref. [10] are depicted in Fig. 2(a) and 2(b), respectively. Predictably, both moments grow up in the *T*-range of 3 K – 120 K and decrease in the range of 150 K – 300 K. The absolue values of the derived orbital moments are quite small, and most importantly, the 3*d* and 4*p* moments gain only factors of ~2 and 3, respectively, as temperature increases from 3 K to 120 K. Magnetic moment obtained by the bulk SQUID-magnetization measurement at 14 Tesla, shown in Fig. 3, increases by factor of 10 as T grows from 3 K – 25 K to 120 K. Recalling that SQUID-magnetization provides the total magnetic moment of Co ions (the sum of orbital and spin moments) the following two conclusions are in order: (i) Major contribution to the total magnetic moment change and 2 times orbital moment change implied by the LS ->HS transition scenario [2] is not observed in the present study, Fig.3 and Fig 2(a). We would like to point out that the IS state implies smaller than HS state 3*d* orbital moment value. The observed less than factor of 2 gain appears to be a good argument for the LS ->IS transition model [1].



**Fig. 2:** Temperature-dependent 3*d* orbital magnetic moment (**a**) as a integrated area from 7709 eV to 7714 eV of the XMCD spectra at the Fig.1 and 4*p* orbital moment (**b**) calculated from Ref.10 per Co ion for LaCoO<sub>3</sub> single crystal.

The previous soft X-ray XMCD data at the Co  $L_{2,3}$ -edges [2], which reflect mainly the surface contribution, do not correlate with our hard X-ray XMCD at the Co *K*-edge (i.e. bulk contribution). The disagreement could probably be attributed to a significant influence of crystal surface (tansile strain, vacansies) [5, 7-9], which results in a stabilization of higher spin states Co<sup>3+</sup>.

## **References**

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**Fig. 3:** SQUID magnetization curves for LaCoO<sub>3</sub> single crystal at different external magnetic fields.