

**Experiment title:**

Pressure-induced spin-state and phase transitions in the mono-LaCoO₃ and multi-valent metastable La_{0.5}Ba_{0.5}CoO_{2.87} with one modified by pulsed huge-field

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HC-5265

Beamline:

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The cobaltites with perovskite structure are systems with fascinating transport and magnetic properties [1-11]. The parent compound LaCoO₃ (LCO) is a paramagnetic (PM) insulator at low temperature (LT). As temperature increases a gradual spin transition (ST) from a low-spin (LS, $t_{2g}^6 e_g^0$) to intermediate (IS, $t_{2g}^5 e_g^1$) or high (HS, $t_{2g}^4 e_g^2$) spin-state occurs, accompanied with insulator to metal transition at ~ 550 K [1-5]. Previous studies showed that chemical substitutions in LCO can induce a long-range antiferromagnetic (AFM) or ferromagnetic (FM) order [2-11]. For example, low doped La_{1-x}Ba_xCoO_{3-d} (LBCO_{3-d}) becomes an AFM insulator [4,7,8,10,11]. One possible explanation for these phenomena is stabilization of the HS state [6, 11]. Surprisingly, the T_c of these compositions increases under applied pressure [7,8,10,11] while the origin of the magnetostructural phase transition is still unclear. In a medium doped oxygen deficit compound La_{0.5}Ba_{0.5}CoO_{2.75} all cobalt atoms are in AFM state whereas La_{0.5}Ba_{0.5}CoO_{2.87} has a mixture FM+AFM phases at low temperatures. It leads to a large magnetic anisotropy inducing a giant magnetoresistance as well as transformation of crystal structure including the separation into two different pseudocubic phases with sp.gr. *Pm-3m* at 50 K and one cubic phase at 300 K (**Fig.1b, c**).

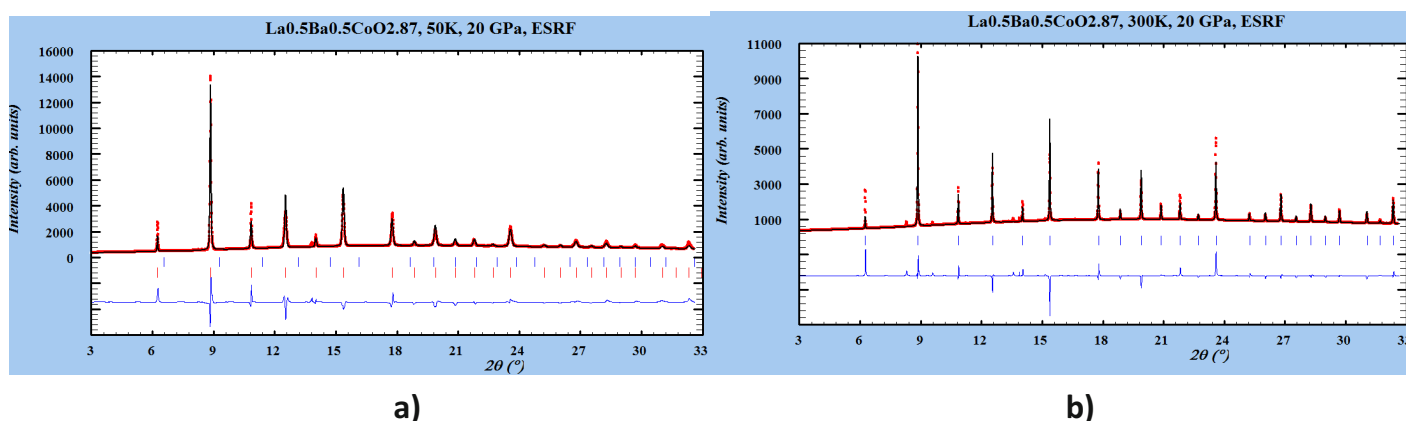


Fig. 1. The Rietveld refinement of the synchrotron X-ray diffraction data for La_{0.5}Ba_{0.5}CoO_{2.87} measured at 20 GPa at 50 K (**a**) and 300 K (**b**), with experimental data in open circles, the calculated pattern in blue, and the difference curve in blue. The tick marks indicate Bragg peak positions for two *Pm-3m* (blue + red) phases at 50 K and one *Pm-3m* (blue) phase at 300 K.

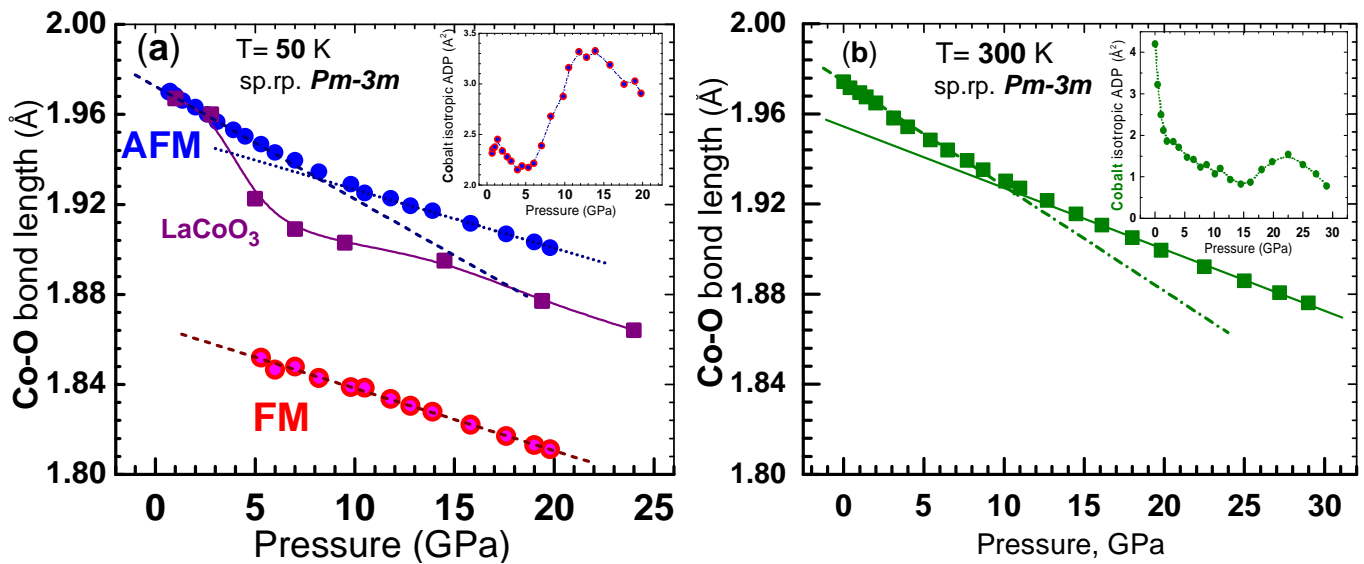


Fig. 2. The pressure dependence of the two (a) and one (b) Co-O bond length for LaCoO₃ with La_{0.5}Ba_{0.5}CoO_{2.87} at 50 K and 300 K, respectively; both insets show an cobalt isotropic ADP.

When pressure is applied, the Co-O bond length of the LaCoO₃ & La_{0.5}Ba_{0.5}CoO_{2.87} is gradually decreasing at 50 K and 300 K (Fig.2a,b). However, above pressure 5 GPa at 50 K for LaCoO₃ & La_{0.5}Ba_{0.5}CoO_{2.87} the Rietveld refinement confirmed a sharply falling of the Co-O bond length and the transformation of one cubic phase into the two ones, respectively (Fig.2a). Moreover, the isotropic ADP of Co for La_{0.5}Ba_{0.5}CoO_{2.87} shows a sharply and weakly growth above 7 GPa at 50 K and above 15 GPa at 300 K, respectively (Fig.2a,b, insets). It interesting to note that these pressure points in both cases corresponds to the bending between higher and lower slopes of Co-O bond length. Similar bending behavior is possibly an indication of the HS→IS spin-state transition of Co³⁺ ions with IS→LS one of Co⁴⁺ at 50 K and IS→LS one of both Co³⁺ and Co⁴⁺ ions at room temperature. The drop of the cobalt isotropic ADP at lower pressures is associated to the lattice compression with further beginning of the HS→IS spin-state transition of Co³⁺ ions and IS→LS one of Co⁴⁺. We suppose that such spin-sate transitions are best described by a mixture of HS, LS and IS multivalent cobalt ions. It is well documented by R.D.Shanon that the Co³⁺ ionic radius of the HS state (0.61 Å) is significantly larger than that of the LS (0.54 Å) and IS (0.56 Å) states and especially with Co⁴⁺ ones of the IS (0.54 Å) and LS (0.53 Å).

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