



$\mu$  the linear absorption coefficient which at this energy  $\mu = 79.3 \text{ cm}^{-1}$ . The sample was carefully cut avoiding as much as possible visible defect creation on the surface. The thickness was 0.14 mm and the measurement of the transmitted beam through the sample yielded an absorption factor of 3.1, which is close to the optimum value. Background was lowered down to 2-3 c/sec with a signal to noise ratio of  $1.5 \cdot 10^4$ , at 4K and at  $Q = (3.5 \ 1 \ 2.5)$

Figure 1 shows the temperature dependence of the peak intensity of the scattering (full circles) and the critical scattering (open circles) in the neighbourhood of the phase transition temperature. Figure 2 shows the extracted hwhm along  $a^*$  (full circles) and  $c^*$  (open circles). Note that the correlation length of fluctuations does not diverge and remains finite at  $T=T_{\text{SP}}$ . There is in addition a second correlation length, shown as full squares in the figure, which has a different temperature dependence. This second length scale is thought to come from clamped order parameter fluctuations in the vicinity of defects.

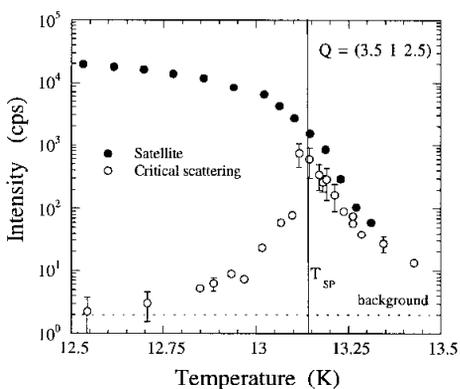


Figure 1.

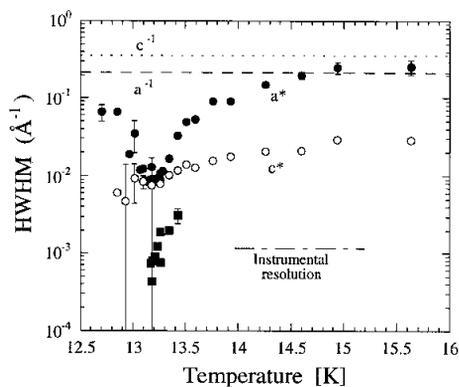


Figure 2.

Not shown in the figures is the anomalous type of line shapes of the critical fluctuations found along both directions. Whereas along  $a^*$  scans have to be fitted with a generalised lorentzian, along  $c^*$  the line shape is gaussian close to  $T_{\text{SP}}$ . The lack of divergence shown in figure 2 and the magnitude of the inverse correlation length along  $a^*$  and  $c^*$  make us think that the atomic displacements giving rise to the superstructure below  $T_{\text{SP}}$  are not the primary order parameter at the phase transition. The phase transition in this compound should be exclusively magnetic in origin, induced by competing antiferromagnetic interactions along the chain axis,  $c$ .